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Citation: Appl. Phys. Lett. 111, 133701 (2017);
View online: https://doi.org/10.1063/1.4998228
View Table of Contents: http://aip.scitation.org/toc/apl/111/13
Published by the American Institute of Physics

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Ferroelectricity and piezoelectricity in soft biological tissue: Porcine aortic walls revisited

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(Received 29 July 2017; accepted 12 September 2017; published online 29 September 2017)

Recently reported piezoresponse force microscopy (PFM) measurements have proposed that porcine aortic walls are ferroelectric. This finding may have great implications for understanding biophysical properties of cardiovascular diseases such as arteriosclerosis. However, the complex anatomical structure of the aortic wall with different extracellular matrices appears unlikely to be ferroelectric. The reason is that a prerequisite for ferroelectricity, which is the spontaneous switching of the polarization, is a polar crystal structure of the material. Although the PFM measurements were performed locally, the phase-voltage hysteresis loops could be reproduced at different positions on the tissue, suggesting that the whole aorta is ferroelectric. To corroborate this hypothesis, we analyzed entire pieces of porcine aorta globally, both with electrical and electromechanical measurements. We show that there is no hysteresis in the electric displacement as well as in the longitudinal strain as a function of applied electric field and that the strain depends on the electric field squared. By using the experimentally determined quasi-static permittivity and Young’s modulus of the fixed aorta, we show that the strain can quantitatively be explained by Maxwell stress and electrostriction, meaning that the aortic wall is neither piezoelectric nor ferroelectric, but behaves as a regular dielectric material. Published by AIP Publishing.

[http://dx.doi.org/10.1063/1.4998228]

Piezoelectricity in bones was first reported in 1957. The main components of bone are the mineral hydroxyapatite \( \text{Ca}_10(\text{PO}_4)_6(\text{OH})_2 \) and collagen. The unit cell of hydroxyapatite, \( \text{P}6_3/m \), is centrosymmetric and, therefore, this natural mineral cannot be piezoelectric. We note that polymorphs with non-centrosymmetric unit cells have been predicted. Recently, piezoelectricity was confirmed in synthetic hydroxyapatite, namely, in electrically poled, textured ceramics and in nanocrystalline thin films. In short, the piezoelectricity in bones is due to collagen. The polypeptide chains in collagen are arranged in a coiled-coil triple helix forming rod-like molecules with short-range crystallinity. The quasi-hexagonal packing of oriented collagen molecules, with symmetry group \( \text{C}_6 \), leads to shear piezoelectricity in the axial direction; dry bone polarizes when a shearing force acts on the oriented collagen fibers so that they slip past one another. The piezoelectric charge coefficient, \( d_{14} \), for dry horse femur and dry bovine bone is in the order of 0.2 pC/N. Other constants, such as \( d_{33} \) in the longitudinal direction, are negligible.

Both the direct and the converse piezoelectric effect have been demonstrated. The extracted linear relationship between electrical displacement and stress, as well as between electric field and strain, proves that bones are truly piezoelectric, as they demonstrate interconversion between electrical charge and mechanical strain. The electromechanical coupling of oriented collagen fibers can have significant consequences for human physiology. For instance, it might explain Wolff’s law, stating that bone in a healthy person or animal remodels itself under an applied mechanical load. We note also that, recently, piezoelectricity has been suggested as a toughening mechanism in seashells.

Piezoelectricity has been reported not only in calcified biological tissue, such as bone and teeth, but also in soft biological tissues such as ligament, trachea, and intestines. The electromechanical coupling is not limited to collagen, but observed in a variety of biopolymers such as cellulose, elastin, keratin, and chitin, which suggests that all fibrous molecules in an oriented state are piezoelectric.

By definition, polycrystalline materials with randomly oriented grains must be ferroelectric in order to exhibit piezoelectric properties. Ferroelectricity means that the material exhibits a spontaneous polarization, which can be switched by an electric field. However, despite the “near-ubiquitous presence of piezoelectricity in biological systems,” there has been very limited evidence for ferroelectricity, especially in soft tissues. The first reports of a ferroelectric response in porcine aortic walls by means of piezoresponse force microscopy (PFM) were therefore surprising. PFM detects the
local deformation of a sample caused by an applied electric field from the tip of the cantilever of a scanning force microscope.\textsuperscript{25,26} For the porcine aorta, the reversal in the piezoresponse phase occurred at about 10 V leading to an estimated coercive field of only 10 kV/m.\textsuperscript{21} Three orders of magnitude smaller than typical ferroelectric polymers such as polyvinylidenefluoride (PVDF).\textsuperscript{27} It was argued that the asymmetry in the coercive voltage could reflect the existence of an internal polarization, as the aorta is internally biased outward.\textsuperscript{28} The phase contrast was approximately 180°, which is a clear indication of polarization switching. Associated with the phase reversal, the deformation-voltage butterfly loops were also observed. The piezoelectric charge coefficient was estimated to be 1 pm/V.\textsuperscript{21} It was noted that this value is two orders of magnitude larger than the one previously reported for blood vessel walls, measured on a macroscopic scale.\textsuperscript{16} Although the PFM measurements were performed locally, the phase-voltage hysteresis loops obtained at different points of the sample were consistent, which indicated global ferroelectricity of the whole aortic tissue.

The ferroelectricity was claimed to be due to the presence of elastin,\textsuperscript{22} an important extracellular matrix protein found in connective tissue. The fibrillar component of elastin consists of simple amino acids such as glycine, valine and alanine.\textsuperscript{29} Both crystalline γ-glycine\textsuperscript{30,31} and β-glycine\textsuperscript{32} are ferroelectric, yet the aorta is not a single crystal. Hence, the reported PFM switching in both aortic walls and elastin was unexpected, and has generated a lot of excitement, but considerable skepticism remains. The latter was expressed, for instance, in the following statement: “Readers should keep in mind that crystalline ferroelectrics are not known to exist in humans, and no in vitro or even in situ work has been reported.”\textsuperscript{33} Here, we contribute to the scientific discussion on biological ferroelectricity by analyzing porcine aorta globally instead of only locally. To that end, we measured both electric displacement and strain as a function of applied electric field on a 1 cm\textsuperscript{2} piece of porcine aorta. We observed neither ferroelectric nor piezoelectric response of the fixated tissue. The strain is proportional to the electric field squared, which indicates that the strain is due to both Maxwell stress and electrostriction, phenomena that occur in any dielectric material.

The aorta specimens were prepared following a procedure reported in the Supporting Information of Ref. 21. In short, about 12 cm of aortas (middle section) was harvested from pigs (\textit{sus scrofa domesticus}), which were part of an experimental lung physiology study at the University Medical Center of Mainz. After removal of surrounding fat and tissue, the aortas were immediately immersed in ice-cold phosphate-buffered saline (PBS). The tissue was always kept at 4°C until final fixation. Figure 1(a) shows a photograph of a piece of the descending thoracic aorta.

For histology, cleaned aortic samples were snap frozen and sectioned into 10 μm slices with a cryostat (HM 560 Cryo-Stat, Thermo Scientific). In order to illustrate cells and elastic fibers, a van Gieson stain was applied. Histological images were acquired using a conventional light microscope (AxioVert200, Zeiss) at 10× and 40× magnification and are shown in Fig. 1(b). The general structure of blood vessels can be subdivided into the intima, media, and adventitia from the inside out.\textsuperscript{34–37} The intima with its endothelial cells delimits the vessel wall from the lumen. The media with smooth muscle cells and connective tissue such as elastin fibers [see inset of Fig. 1(b)] accounts for the majority of the vascular wall, especially in big arteries. The adventitia represents the border to the tissue surrounding the vessel, frequently with its own blood supply.

For electrical and mechanical characterization, porcine aortic samples with an area of about 1 cm\textsuperscript{2} and a thickness of 1 mm were fixated using 4% paraformaldehyde in PBS for 1 h at ambient temperature. Then, samples were dehydrated in an ascending ethanol/deionized water series for 15 min each (30/70; 50/50; 70/30; 100/0). Afterwards samples were subjected to an ascending hexamethyldisilazane/ethanol series (30/70; 50/50; 70/30; 100/0) for 15 min each and finally dried overnight in a hood.

Uniaxial tensile tests were performed on an Instron Model 3365 universal testing system equipped with a 1 kN load cell. Tensile specimens of fixated porcine aorta, having an approximate width of 5 mm, a thickness of 1 mm and a length of 4 mm, were stretched between the clamps with a tensile rate of 1 mm/min at ambient temperature.

The electric displacement as a function of electric field was measured using a Radiant precision multiferroic test system (Radiant Technologies, Inc.) at a frequency of 10 Hz. Simultaneously, the strain as a function of electric field was measured using a MTI 2100 photonic sensor interfaced with the Radiant tester. Piezoelectricity was additionally investigated using a Berlincourt-type piezometer (PM300, Piezotest, London, UK). A static force of 10 N was used under a 0.25 N peak to peak sinusoidal excitation at 110 Hz. Impedance was measured using a Schlumberger Si 1260 Impedance Analyzer. All electrical measurements were performed in ambient conditions. We emphasize that we tested aortas from two different pigs and various pieces from each aorta. Our results could be reproduced in all the tests.

Figure 2(a) shows the electric displacement versus electric field. Electrical conduction can be excluded as the dielectric loss is less than 1%. A linear relation is obtained up to the maximum electric field of about 6 MV/m, which is
where the aorta, linear contribution to the displacement can be disregarded for on the electric field and the poling history. However, this non-aligned electric dipoles in a ferroelectric material and depends dielectric constant. The spontaneous polarization is due to electric field, as presented in Fig.2(b). The strain is negative; porcine aortic wall, we measured the strain as a function of charging of any linear dielectric capacitor. The aortic specimen behaves as a normal dielectric.

The electrical displacement, \( D \), for a capacitor under an applied electric field, \( E \), is the sum of the induced polarization, \( P'_r \), and spontaneous polarization \( P_s \):

\[
D = P_i + P_s = \varepsilon_0 \varepsilon_r E + P_s = \varepsilon_0 \varepsilon_r E, \tag{1}
\]

where \( \varepsilon_0 \) is the vacuum permittivity and \( \varepsilon_r \) is the relative dielectric constant. The spontaneous polarization is due to aligned electric dipoles in a ferroelectric material and depends on the electric field and the poling history. However, this non-linear contribution to the displacement can be disregarded for the aorta, \( P_s = 0 \), as there is no hysteresis in the displacement as function of electric field, cf. Fig. 2(a). There is no sign of a ferroelectric spontaneous polarization. The displacement is only due to the induced polarization, which represents the charging of any linear dielectric capacitor.

To investigate the electromechanical properties of the porcine aortic wall, we measured the strain as a function of electric field, as presented in Fig. 2(b). The strain is negative; the aorta specimen contracts when an electric field is applied. Here again, there is no indication of hysteresis; a butterfly shape as expected for ferroelectric materials is not observed. The shape of the strain versus electrical displacement curve looks like a parabola. As verification we present the strain as a function of electric field squared in Fig. 2(c). A straight line is indeed obtained and the least square approximation yields a slope of \(-2.08 \times 10^{-18} \text{ m}^2/\text{V}^2\). This value also allows fitting the measured data in Fig. 2(b), as indicated by the red lines.

The aortic wall is hence not ferroelectric, but also not piezoelectric as additionally supported by Berlincourt-type piezometer measurements. No measurable signal could be detected for \( d_{33}, d_{33} \), and \( d_{15} \) demonstrating that the aortic wall is not piezoelectric, i.e., there is no surface charge generated under mechanical stress.

The strain induced in a non-piezoelectric, isotropic dielectric material by an electrostatic field has two sources. The first one is the Maxwell strain, which is due to the electrostatic force resulting from the free charges on the electrodes of a capacitor. The second one, referred to as electrostriction, is a universal property of solid and liquid dielectrics arising from changes in the dielectric constant with the electrically induced strain. The total longitudinal strain, \( S_{33} \), is given by

\[
S_{33} = -\frac{1}{2} \varepsilon_0 \varepsilon_r E^2/Y \times \left( 1 + 2\nu \right) \left( a_1 + a_2 (1 - 2\nu) \right), \tag{2}
\]

where \( Y \) is the Young’s modulus, \( \nu \) is the Poisson ratio, and \( a_1 \) and \( a_2 \) are two electrostrictive parameters describing the change in dielectric properties of the material under shear and bulk deformation, respectively. This equation holds for non-compliant electrodes; for compliant electrodes, the longitudinal strain is twice as large. When the dielectric constant is independent of deformation, then \( a_1 \) and \( a_2 \) vanish and only the compressive Maxwell strain remains, given by the first term of Eq. (2) as

\[
S_{33}\text{Maxwell} = -\frac{1}{2} \varepsilon_0 \varepsilon_r E^2/Y \times (1 + 2\nu), \tag{3}
\]

and the pure electrostrictive contribution is given by the second term of Eq. (2) as...
where $Q_{33}$ is the commonly used longitudinal electrostrictive coefficient, derived from the phenomenological Devonshire theory.\(^{40}\) For ionic insulators, the electrostrictive strain is dominated by anharmonicity in the electrostatic potential leading to a positive $Q_{33}$.\(^{41}\) For polymers, covalent bonding is much stronger than intermolecular forces such as van der Waals interactions. This effectively leads to a negative value of $Q_{33}$, and, hence, to a compressive strain. Both the Maxwell- and electrostrictive strain depend on the squared electric field and are therefore indistinguishable.

In order to quantitatively analyze the strain, a value for the Young’s modulus is required. Since the porcine aorta samples underwent fixation before electrical characterization, we cannot use literature values reported for fresh aortas. Therefore, we performed a tensile test for a piece of fixated porcine aorta. A typical stress-strain curve in the linear regime is presented in Fig. 2(d). From the slope between 1.0% and 2.0% strain, we extract an average Young’s modulus of 160 ± 20 MPa. The process of fixation increases the stiffness of the aorta piece by two orders of magnitude with respect to a fresh porcine aorta due to denaturation of the proteins by alcohol and cross-linking by paraformaldehyde.

We calculate from Eq. (3) the Maxwell strain by taking the extracted value of 26 for the static dielectric constant, and a Poisson ratio of 0.3.\(^{45}\) We calculate for the proportionality constant between strain and electric field squared a value of $-1.1 \times 10^{-18}$ m$^2$/V$^2$. This value is about half of the experimentally determined slope, cf. Fig. 2(c), meaning that about half of the strain is due to Maxwell stress. The other half of the strain is due to electrostriction, leading to an extracted electrostrictive coefficient, $Q_{33}$, of $-17.4$ m$^2$/C$^2$. We note that this coefficient can be estimated from the empirical linear relation for the absolute value of the hydrostatic electrostrictive constant $Q_3 = 1/Y_0 e_0 Y$.\(^{40,46}\) A theoretical derivation for electrostrictive polymers is reported based on a microscopic Debye/Langevin formalism.\(^{47}\) The calculated value of 27 m$^3$/C$^2$ is in fair agreement with the experimentally extracted value for the aortic wall.

Our measurements at electric fields up to 6 MV/m allowed us to demonstrate that the strain as a function of electric field is a parabola, i.e., the strain depends on the electric field squared. In contrast, reported PFM measurements were performed at much lower fields up to only 0.1 MV/m. To compare the local and global datasets, we approximated our global strain/electric field curve as a straight line. At low fields, we then obtain a slope of 2 pm/V, which could explain the value of 1 pm/V as extracted from the reported local PFM measurements.

Our global strain measurements on aortic walls are at variance with reported PFM measurements.\(^{21-24}\) Piezoresponse force microscopy is a powerful tool and well-established in the ferroelectric community.\(^{25,26,48-51}\) However, it was demonstrated to be prone to artifacts.\(^{52-54}\) We deliberately tried, but failed to reproduce the PFM phase switching. The reported switching might be due to voltage gated ion channels or reorientation of dipoles, which are “internally biased outward” in the aorta. These responses are not ferroelectric but remain intriguing and open for further investigation.

We gratefully acknowledge technical support from C. Bauer, H. Raich, and F. Keller, and stimulating discussions with P. W. M. Blom, all from the Max Plank Institute for Polymer Research, Mainz, Germany. We want to thank C. Möllmann from the University Medical Center for help with the extraction of the aorta. T.L. acknowledges financial support by the Graduate School Materials Science in Mainz.

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