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Geometry of the Magma Chamber and Curie Point Depth Beneath Hawaii Island: Inferences From Magnetic and Gravity Data

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This study used land gravity and airborne magnetic data to investigate the depth to the magmatic chamber and map the heat flow distribution beneath the active volcanoes of Hawaii Island using the Curie point depth (CPD) and gravity modeling. Obtaining some of the ground-based geophysical measurements was problematic due to accessibility limitations; therefore, this study used available data. The CPD and magnetic data were used to map the depth to the bottom of the magnetic layer by calculating the depth to the Curie isotherm (540°C) beneath Hawaii Island. The spectral peak method was used to calculate the depths to the shallow and deep magnetic sources for the entire island, and the CPD was calculated using the centroid method. A two-dimensional density model for two Earth layers was constructed using forward modeling of the gravity data. A large plume of dense intrusive material was observed beneath the three adjacent volcanoes of Mauna Loa, Mauna Kea, and Kilauea, and two small chambers were found to be located beneath the Kohala and Hualalai volcanoes. Based on the gravity modeling results, the depth to the magma layer varied from 0.5 to 10 km, and the heat flow was higher close to the volcanic eruption zones. The current study is informative and cost effective for the world's most active volcanic areas.

Keywords: magmatic chamber, gravity inversion, Curie depth point, geothermal gradient, heat flow, Hawaii Island

INTRODUCTION

The structural complexity of volcanic islands and their difficult accessibility make geophysical imaging problematic. Moreover, geophysical modeling generally suffers from nonuniqueness in its solutions. This study uses two geophysical techniques to overcome these problems and reduce modeling errors. Despite the accessibility limitations, geophysical modeling techniques of potential data have been used extensively in different geologic settings to provide reliable results on structures and geometries (Green, 1975; Guillen and Menichetti, 1984; Camacho et al., 1997; Li and Oldenburg, 1998; Camacho et al., 2000; Boulanger and Chouteau, 2001; Montesinos et al., 2006; Cella et al., 2007; Camacho et al., 2011; Marcotte et al., 2014; Barnoud et al., 2016). **Figure 1** shows the location and topography of the Hawaiian Island arch system.

Geophysical and structural measurements are used extensively to delineate the three-dimensional (3-D) shape of granitoid plutons (Miller and Tauch, 1989; Vigneresse, 1995a, b; Eguiluz et al., 1999)

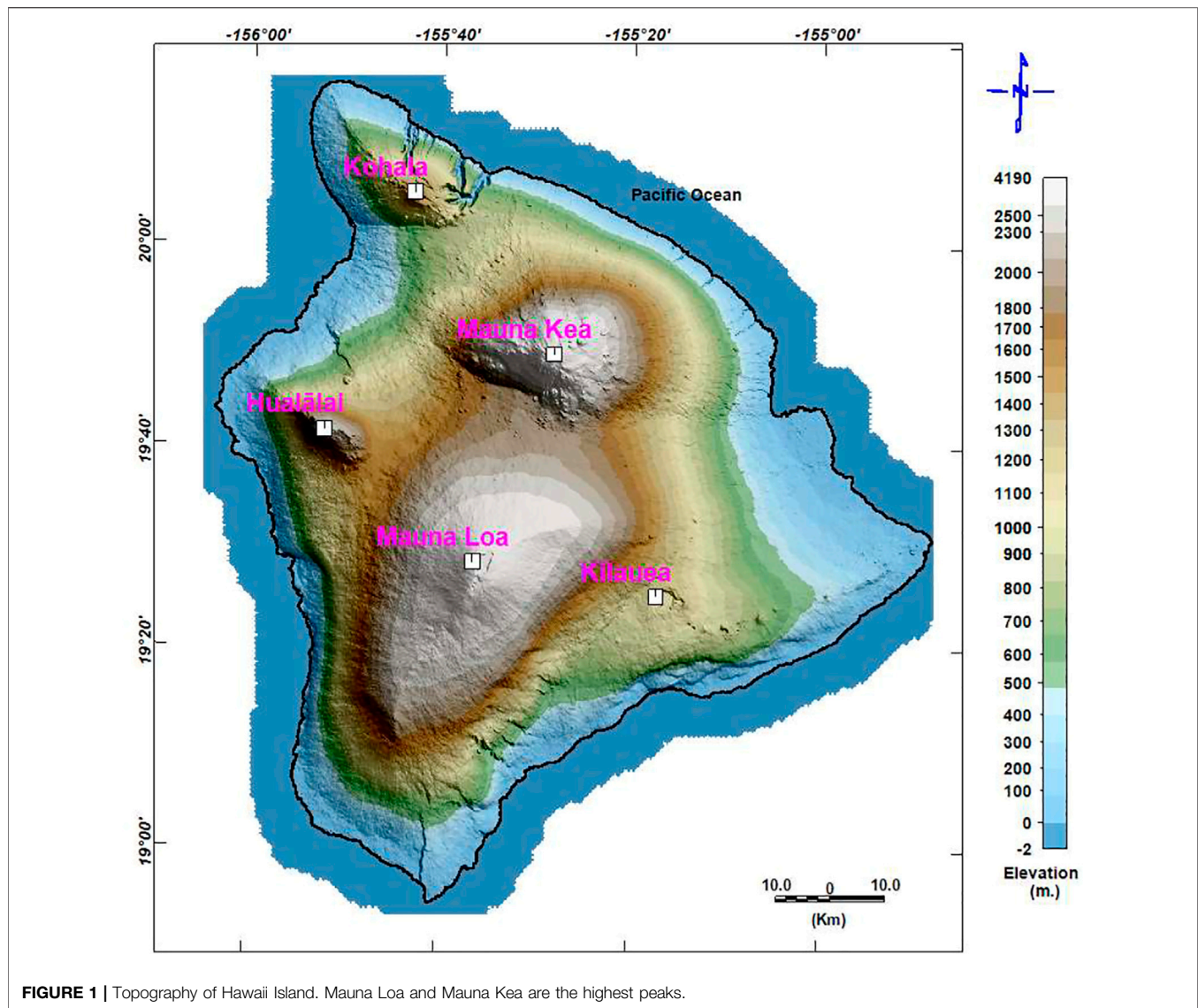


FIGURE 1 | Topography of Hawaii Island. Mauna Loa and Mauna Kea are the highest peaks.

in different areas. Most previous studies include horizontal map views of intrusions in their results. However, scientific studies assume the widespread concept of the inverted teardrop when considering pluton 3-D geometries (Vigneresse, 1995b) based on structural extrapolations to depth. Most field-based studies use extrapolations to depth without direct evidence to obtain a reasonable approximation when geophysical techniques are employed.

Aeromagnetic datasets are widely used to estimate the Curie point depth (CPD) and assess geothermal resources (Chiozzi et al., 2005; Trifonova et al., 2009; Hsieh et al., 2014). The CPD is considered at the surface at a high temperature of $\sim 540^{\circ}\text{C}$, as measured *in situ* from holes drilled through the crust and into the still-molten lens of tholeiitic basalt in Kilauea (Zablocki and Tilling, 1976). The CPD surface indicates the bottom of the magnetic layer; this is where paramagnetic minerals are formed from ferromagnetic minerals when they reach the Curie point temperature (Nagata, 1961) or the depth at which

nonmagnetic rocks are formed from magnetic rocks (Ravat et al., 2007). Iron sulfides and iron-titanium oxide are the most common ferromagnetic minerals (Rajaram, 2007).

Castro and Brown (1987) have used high-resolution sampling of the 1950 and 1972 flows in Kilauea to identify the intraflow variations of very young basalt flows. The paleomagnetic directions of these samples are different from the direction of Hawaii's geomagnetic field. This shallow inclination anomaly was attributed to the recording mechanism of subaerial basalts (Castro and Brown, 1987); however, it has since been identified as the nondipole field effect (Cox, 1975). Another paleomagnetic study has identified the lavas in the main part of the Pohue Bay flow and those of the Hawaiian cones to have similar paleomagnetic properties with relatively similar ages (Jurado-Chichay et al., 1993). That study could have identified both younger lava flows, which later utilized the main tube, and an earlier subset of cones, which were formed before the Pohue Bay eruption.

Magnetic data have been analyzed extensively to estimate both the depth to the magnetic basement (Spector and Grant, 1970; Hahn et al., 1976; Garcia-Abdeslem and Ness, 1994) and the Curie isotherm (Blakely, 1988; Okubo and Matsunaga, 1994). According to Spector and Grant (1970), the depth factor controls the shape of the readily averaged power spectrum. Their statement has resulted in a wide interpretation of the average two-dimensional (2-D) power spectrum of magnetic data. Moreover, the depth to the source can be calculated from the slope of the log radially averaged power spectrum.

The thermal structure of the crust in different geologic environments has been widely investigated using magnetic data (Spector and Grant, 1970; Bhattacharyya B. K. and Leu L.-K., 1975, 1977; Byerly and Stolt, 1977; Okubo et al., 1985; Blakely, 1995; Tanaka et al., 1999; Chiozzi et al., 2005; Ross et al., 2006; Trifonova et al., 2009; Gabriel et al., 2011, 2012; Bansal et al., 2013, 2016; Hsieh et al., 2014; Nwankwo and Shehu, 2015; Nwankwo and Abayomi, 2017); the resulting geomagnetic anomalies above the CPD have been used to delineate magnetic structures (Bhattacharyya and Leu, 1975a, b; Byerly and Stolt, 1977; Blakely and Hassanzadeh, 1981; Blakely, 1988; Chiozzi et al., 2005; Trifonova et al., 2009; Hsieh et al., 2014; Idárraga-García and Vargas, 2018; Mohamed Al Deep, 2021).

Flinders et al. (2013) used land and marine gravity data and a 3-D gravity model to calculate the average densities, volumes, and percentages of olivine in the intrusive materials and cumulate cores below the volcanoes in Hawaii. They used an isosurface density of 2.85 g cm^{-3} to delineate intrusive material, which equated to over 60% of dikes with a density of 2.95 g cm^{-3} . Furthermore, they defined cumulate cores using an isosurface density of 3.00 g cm^{-3} , which corresponded to ~35% olivine (density: $3.2\text{--}3.3 \text{ g cm}^{-3}$) in the intrusive complex (density: 2.85 g cm^{-3}).

The Hawaiian Islands were formed during the past 70 Ma as the Pacific lithospheric plate moved north and then west relative to a melting anomaly. This phenomenon is represented by the hotspot hypothesis, which accounts for the formation of a volcanic chain on the ocean floor. Wilson (1963a, c) proposed that the islands of Hawaii were formed when the seafloor moved over lava sources in the asthenosphere; they limited this theory to the volcanoes and the ridge of the Hawaiian Islands. Subsequently, Wilson's theory was expanded by Christofferson (1968) to include the Emperor Seamounts. Morgan (1972a, b) suggested that the Hawaiian hotspot and others are thermal plumes of material rising from the deep mantle. Clague and Dalrymple (1987) tested and validated this hypothesis based on a study of the geologic evolution of the Hawaiian–Emperor volcanic chain. In terms of geology, tholeiitic basalts account for >95% of the islands' rocks (Clague and Dalrymple, 1989).

Geophysical data collected from the ground and from the airborne were commonly employed in groundwater investigations and subsurface geology (e.g., Meneisy and Al Deep, 2020; Mohamed and Abu El Ella, 2021; Al Deep et al., 2021). On the other hand, Global data from the Earth Gravitational Model and the Earth Magnetic Anomaly Grid have been widely used for crustal studies and the depth to the bottom of the magnetic layer (e.g., Idárraga-García and Vargas,

2018; Mohamed and Al Deep, 2021), whereas gravity data from the Gravity Recovery and Climate Experiment mission have been successfully applied for estimating mass transport and distribution in the Earth's fluid (e.g., Mohamed et al., 2017; Mohamed, 2019; Mohamed, 2020a; Mohamed, 2020b; Mohamed, 2020c; Taha et al., 2021; Mohamed and Gonçalves, 2021; Mohamed et al., 2021; Mohamed et al., 2022).

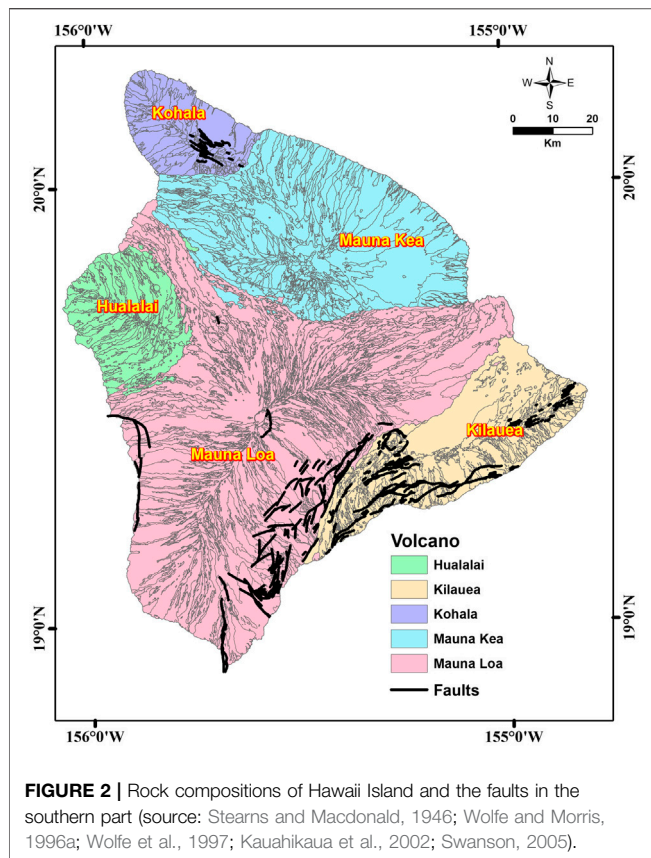
The current study aims to characterize the morphology of the magmatic chamber beneath Hawaii Island using gravity and magnetic potential fields. Delineating the geometrical border of the magmatic chamber beneath the island of Hawaii and estimating the CPD, geothermal gradient, and heat flow of the area require the techniques of gravity inversion and spectral analysis of magnetic data. Our method of estimating the CPD to determine the maximum depth to the magnetized rocks calculates the depth to the upper surface of the magma chamber indirectly. Additionally, a two-layer 2-D geological model using the contrast in density between the average density of the solidified rocks and the magma layer is applied using gravity modeling techniques.

GEOLOGIC SETTING OF HAWAII ISLAND

Hawaii Island comprises five major shield volcanoes (**Figure 2**): Mauna Kea (MK), Kohala (Ko), Kilauea (Ki), Mauna Loa (ML), and Hualalai (Hu) (Langenheim and Clague, 1987). Except for some minimal erosion to the northern sides of Ko and MK, these volcanoes have suffered little erosion (Langenheim and Clague, 1987). They progress in age from the southeast end (with its still-active volcanoes) to the northwest end (with its volcanoes dating from 75 to 80 Ma) (Clague and Dalrymple, 1987).

The ML volcano is the largest on Earth at 97-km long, 48-km wide, and an estimated volume of 70–80 million km^3 (Robinson and Eakins, 2006). Its shield-shaped dome rises to 4,167 m above sea level (**Figures 1, 2**), and its slopes are dotted with a few cinder cones. The caldera of Mokuaweoweo is located on the summit (Macdonald, 1977). ML formed over a period of ~0.5 million year. During recent years, lava younger than 1,000 year has poured from fissures and a cone on the volcano's floor and rim, filling depressions and covering about 40% of its surface area (Lockwood and Lipman, 1987). A few weak explosions have occurred. The surface of ML originates mostly from the Holocene (Stearns and Macdonald, 1946; Lipman and Swenson, 1984). Stearns and Macdonald (1946) divided the volcano's rocks into Pliocene Ninole basalt, Pleistocene Kahuku basalt, and Pleistocene and Holocene Kau basalt (the most recent).

At 4,205 m above sea level, MK is Hawaii's highest volcano (**Figures 1, 2**). There have been no eruptions over the past 3,600 years (Porter, 1979a). The configurations of cinder cones indicate the location of less-defined southerly, easterly, and westerly rifts. The volcano's rocks were divided into two volcanic series by Stearns and Macdonald (1946). The older (Hamakua) series forms a major part of the mountain and represents the shield stage and part of the post-shield stage. The upper part of the mountain (above 3,353 m) comprises a plateau resulting from the Laupahoehoe series filling of a caldera



in the Hamakua volcanic series, which represents the rest of the post-shield stage. Porter (1979a, 1979b) redefined the volcano's two rocks to include glacial deposits; additionally, they were elevated to a group ranking of volcanic and glacial formations.

The last eruptions of the Hu volcano in 1800–1801 occurred from five separate vents (Moore et al., 1987) (Figures 1, 2). The volcano has an approximate diameter of 27 km and a well-defined NW-striking rift zone; its less-well-defined N- and SE-trending rift zones are characterized by widespread cinders and spatter cones (Langenheim and Clague, 1987). Its subaerial surface is represented by the Pleistocene Waawaa Trachyte Member, with short flows of hawaiite and a few flows of alkali basalt lava (Moore et al., 1987). According to detailed mapping and C-14 dating, the volcano's Holocene lava flows have been divided into four groups: 10–5, 5–3, 3–1, and <1 ka (Moore et al., 1987).

Based on K–Ar analysis, Ko is the oldest Pleistocene (McDougall and Swanson, 1972) and longest inactive volcano in the islands (Lipman, 1980) (Figures 1, 2). It was formed over northwesterly and southwesterly rifts and a weak southwesterly rift (Langenheim and Clague, 1987). Its last eruption was ~60,000 years ago, and it is considered to be extinct. The volcano's rocks represent two volcanic series (Stearns and Macdonald, 1946): the older Pololu series (Pololu basalt) (Langenheim and Clague, 1987) of shield-stage tholeiitic basalt and caldera-filling postshield-stage alkalic basalt (Stearns and Macdonald, 1946) and the younger Hawi series (Hawi Volcanics) (Langenheim and

Clague, 1987) of differentiated alkalic lava of the postshield stage (Stearns and Macdonald, 1946).

The Ki volcano has an area of 2,500 km² and rests on the southeast slope of ML (Figures 1, 2). Lava flows from ML pass over the slopes of Ki. The Ki volcano is the youngest in Hawaii and remains very active (Clague and Dalrymple, 1987). Around 70% of its surface is younger than 500 years, and about 90% is younger than 1,100 years (Holcomb, 1987). At the base of a fault escarpment on its mobile south flank, Ki's older rocks of Hilina basalt are between 100 and 30 kyr old (Easton, 1987). The most recent Puna basalts of the Pleistocene and Holocene (Easton, 1987) represent the younger rocks and are separated from the older ones by Pahala ash.

DATA

The land and airborne potential field data used in this study are described below.

Gravity anomaly data

Gravity measurements from the study area were collected as datasets on the deep magmatic structures of the Hawaiian volcanoes. A complete Bouguer anomaly map was constructed from the corrected observed gravity data to build a realistic model for Hawaii relative to the ground surface. The gravity data, which were collected by Kauahikaua (2017), did not contain base station time series records, and there was no time channel for the data points. Therefore, it was assumed that the instrumental drift had been completed. A free air anomaly map was constructed, and a complete Bouguer anomaly map was compiled for use in the gravity inversion calculation.

The first step of the data processing involved calculating the theoretical gravity for the study area by Eq. 1. Provided the data projection was in the WGS84, the same datum was used.

$$G_n = a_1 \frac{(1 + a_2 (\sin \phi)^2)}{\sqrt{(1 + a_3 (\sin \phi)^2)}} \quad (1)$$

The subtraction of theoretical gravity from the observed data resulted in latitude-corrected data. The formula below (Eq. 2) was used for the latitude correction calculations:

$$G_l = G_{observed} - G_n \quad (2)$$

where G_n is the theoretical gravity, ϕ is the latitude, and a_1 , a_2 , and a_3 are constants that are equal to 9780326.7714, 0.00193185138639, and -0.00669437999013, respectively.

The second step involved the calculation of the free air anomaly (G_{fair}) assuming that the elevation of the ground above or below the datum was without any representation of the density variation. Therefore, this was excluded from this study's modeling. The free air correction was calculated and subtracted from the latitude correction for the WGS84 data using the following equations:

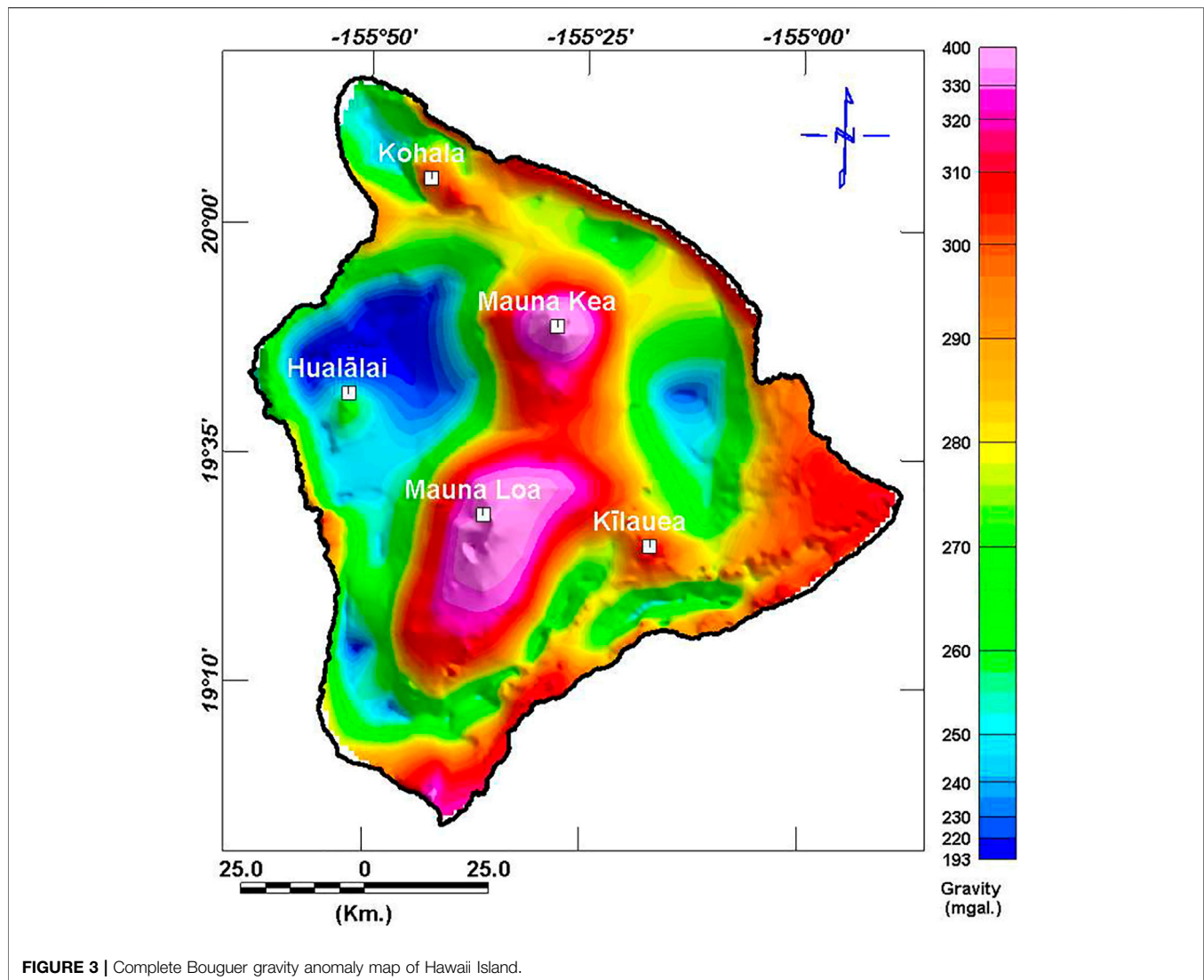


FIGURE 3 | Complete Bouguer gravity anomaly map of Hawaii Island.

$$\delta gh = -(3.083293357 + 0.004397732 * \cos^2 \phi) * h + 7.2125 * 10^{-7} * h^2 \quad (3)$$

$$G_{fair} = G_n - \delta gh \quad (4)$$

where δgh is the free air correction, h is the elevation of the gravity sensor above the datum, and ϕ is the latitude.

A simple Bouguer anomaly map was constructed by subtracting the Bouguer correction, which replaces the air above a datum by the mean density of land. The complete Bouguer correction calculates the terrain correction for the gravity stations and then applies the complete Bouguer correction to the simple Bouguer-corrected data. The terrain correction used topography data obtained from the Shuttle Radar Topography Mission (NASA, 2013).

The final complete Bouguer anomaly (Figure 3) was calculated according to Eq. 5:

$$G_B = G_{fair} - 0.0419088 * [D * H_s + (D_w - D) * H_w + (D_i - D_w) * H_i] - G_c \quad (5)$$

where G_B is the complete Bouguer anomaly, D is the Bouguer density of the Earth in that area ($\sim 2.73 \text{ g cm}^{-3}$), H_s is the station elevation (meters), D_w is the Bouguer density of water ($\sim 1.027 \text{ g cm}^{-3}$), H_w is the water depth (meters), D_i is the Bouguer density of ice ($\sim 0.917 \text{ g cm}^{-3}$), H_i is the ice thickness (meters), and G_c is the correction of the Earth's curvature.

Magnetic anomaly data

The available magnetic data for Hawaii Island comprise four aeromagnetic datasets. The United States Geological Survey's Hawaii-78-Hawaii data are the only data available with complete coverage of the island; they were collected along the north-south lines by aircraft instruments at a height of 305 m, and magnetic field values and locations were recorded (Godson et al., 1981). The data were collected along longitudinal lines with a spacing of 1.6 km and include latitude, longitude, altitude, and magnetic field values. Some variations in magnetic measurements are caused by rocks with high proportions of magnetic minerals; these anomalies reflect variations in the amount/type of magnetic material and the

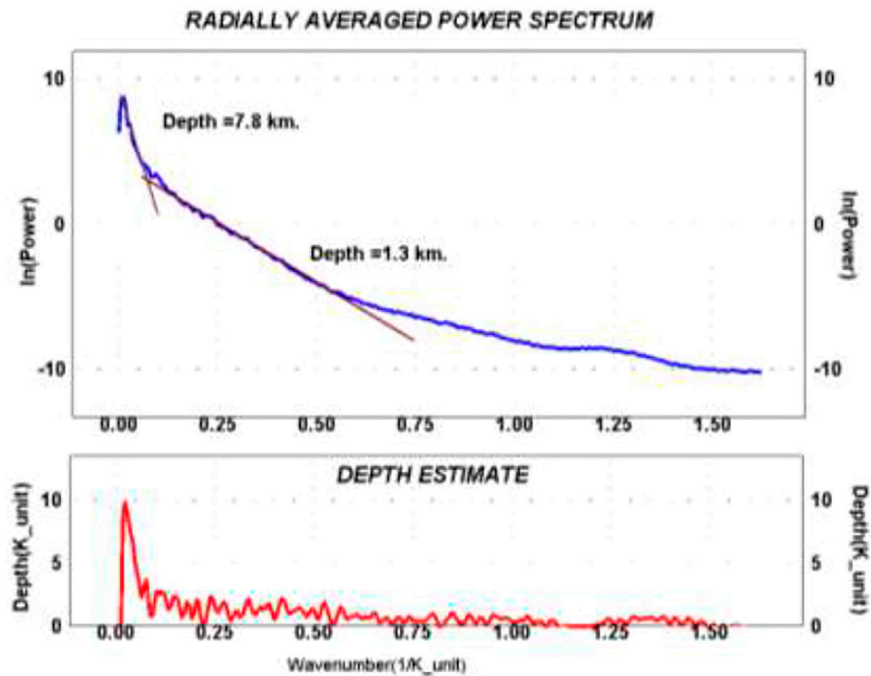


FIGURE 4 | Calculated depth from the averaged 2-D power spectrum to magnetic sources that may relate to volcanic cycles with high magnetite content.

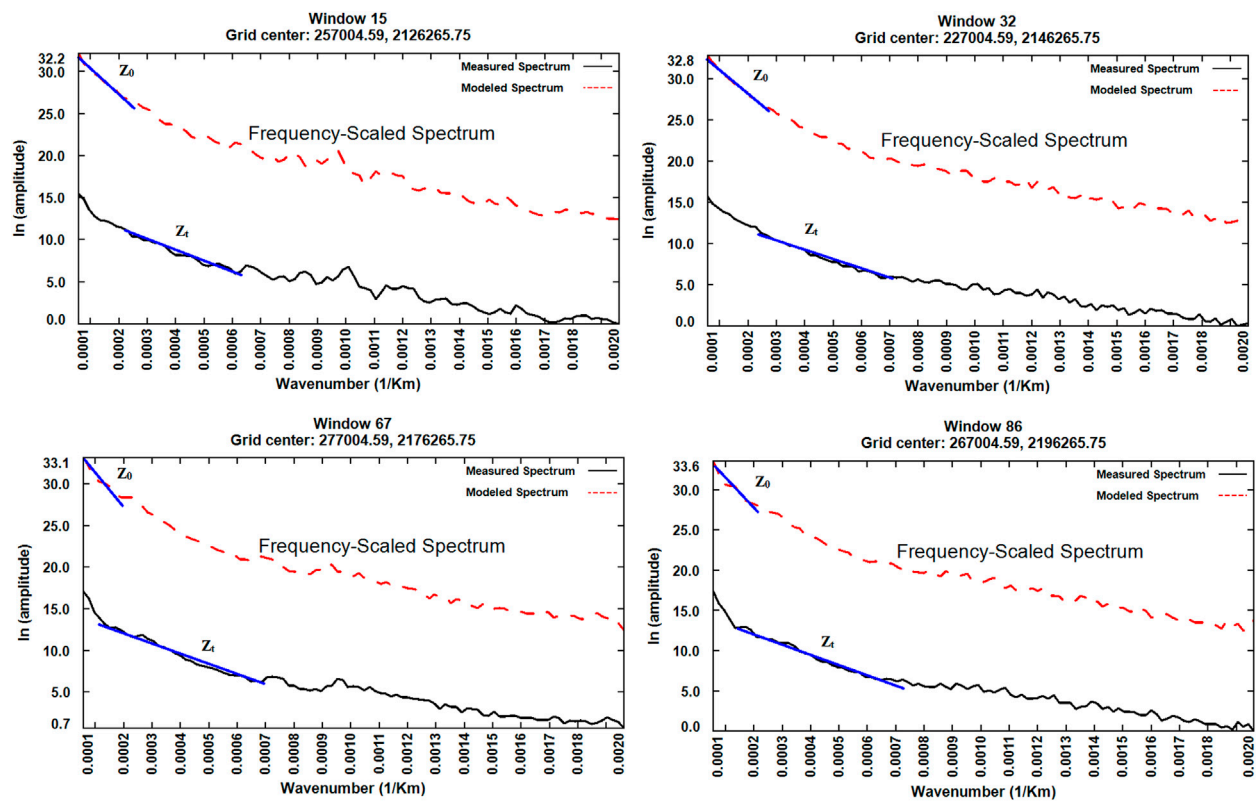


FIGURE 5 | Examples of spectra for the estimation of the depth of the top bound (Z_1) and the depth of the centroid (Z_0) of magnetic sources for different windows.

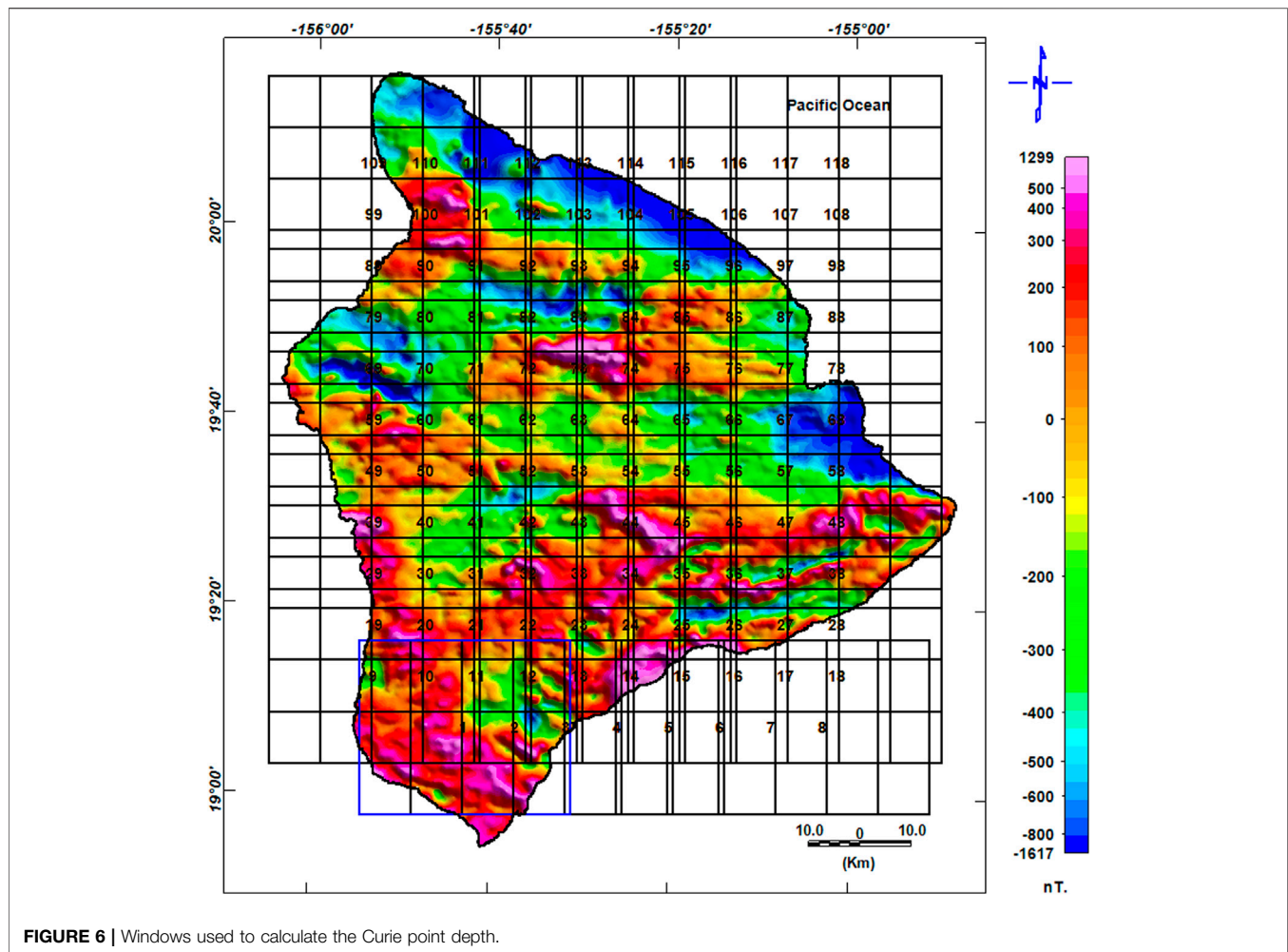


FIGURE 6 | Windows used to calculate the Curie point depth.

shape/depth of the rock bodies. The features and patterns of the aeromagnetic anomalies in the Hawaii-78-Hawaii dataset were processed to map the CPD and estimate the geothermal activity in the active volcanic zones based on the high resolution and consistency of the data along the flight lines.

Hildenbrand et al. (1993) used this dataset as well as other data that were obtained along lines normal to the rifts of ML and Ki at a flight height of 90 m (Flanigan et al., 1986). By merging the two datasets, they produced the most comprehensive dataset for the entire island. This is the only difference in our data reduction, as we used aeromagnetic data that provided complete coverage (Hawaii-78-Hawaii) of the island.

The complete dataset was gridded with 0.5-km spacing using the minimum curvature technique, and the dataset was continued upward to 305 m using the Geosoft 8.4 software (Geosoft Oasis Montaj, 2015). The shifted anomalies were adjusted using a reduction-to-pole transformation, while spectral and centroid techniques were used to delineate the shallow and deep sources and define the CPD, which denotes the interface between the magma layer and the overlying magmatic layer. The terrain effect was calculated using a spectral inversion process, which was subsequently removed to estimate the terrain-corrected magnetic field used in the current study.

RESULTS

Identifying the upper surface of the magmatic chamber from magnetic data

Calculating the CPD to determine the maximum depth to the magnetized rocks allows the upper surface of the magma chamber to be delineated indirectly. The rocks in Iki lava lake in Hawaii cannot contain any magnetization above the Curie temperature (540°C) (Zablocki and Tilling, 1976). Thus, for magnetite (the most abundant magnetic mineral), residual magnetism appears below the Curie point of ~540°C. In this case, the magma beneath Hawaii Island is basic, reflecting higher temperatures of 1,100°C to 1,300°C. The CPD denotes the depth to the bottom of the magnetized rocks, i.e., the depth to the upper surface of the magmatic chamber. As the CPD calculation involves the power spectrum, we calculated the 2-D radial power spectrum for the entire area to provide an overview of the depth distribution in that vicinity.

The spectral peak and centroid methods are commonly used spectral techniques for estimating the depth to the bottom of a magnetic layer (Ravat et al., 2007). The spectral peak method (Spector and Grant, 1970) was used by Shuey et al. (1977) and Connard et al. (1983). The present study also used this method to

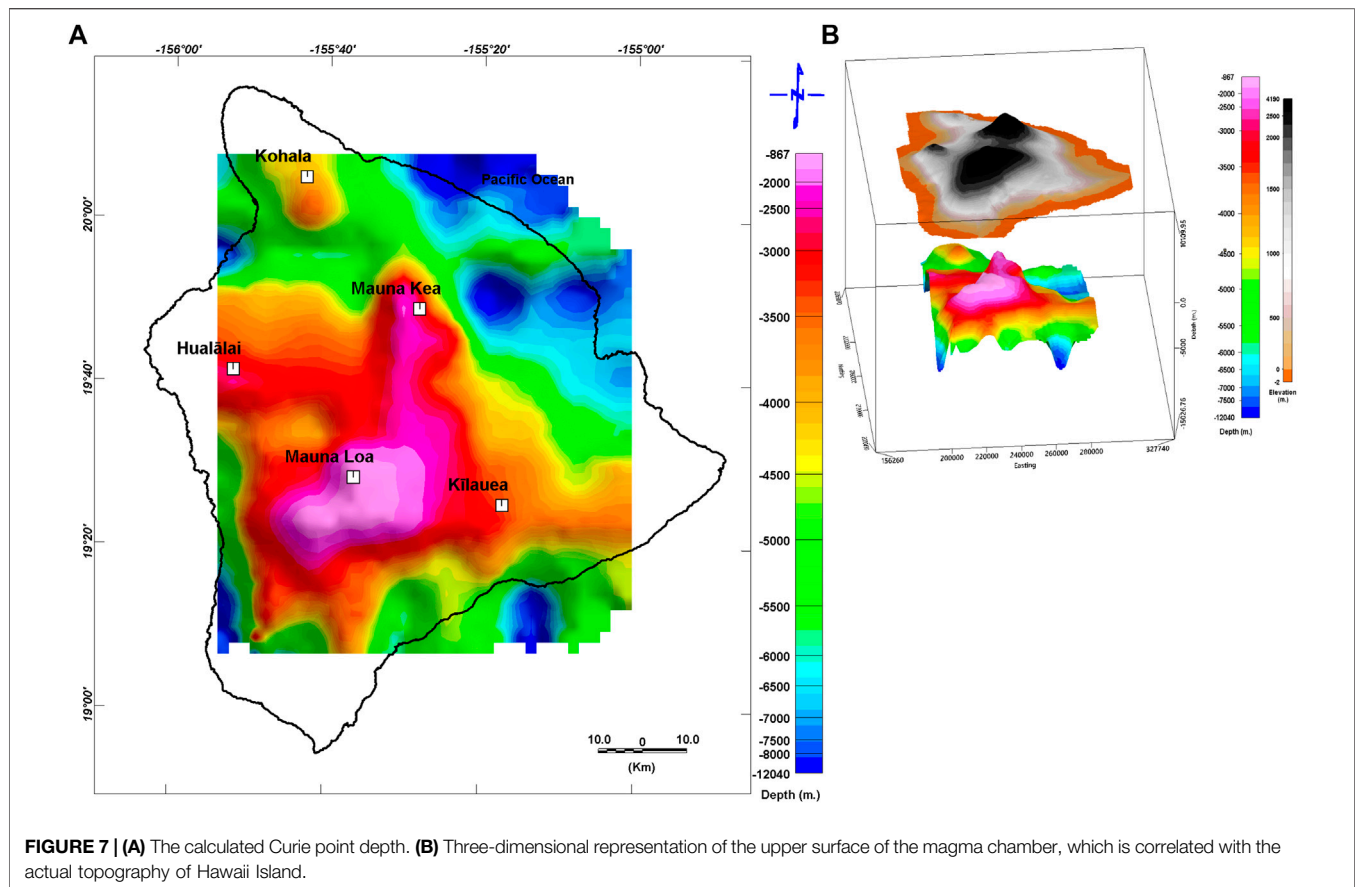


FIGURE 7 | (A) The calculated Curie point depth. **(B)** Three-dimensional representation of the upper surface of the magma chamber, which is correlated with the actual topography of Hawaii Island.

allocate depth to the top of the magnetized sources along with the centroid method (Bhattacharyya B. K. and Leu L. K., 1975, 1977; Okubo et al., 1985; Tanaka et al., 1999) to determine the centroid of rectangular parallelepiped sources or the depth to the centroid.

Spectral peak method (two-dimensional power spectrum)

The Spector and Grant (1970) equation (after Blakely, 1995) was used to estimate the depths to the bottom and top of the collective magnetic sources from their averaged power spectra:

$$|F(k)|^2 = 4\pi^2 C_m^2 |\theta_m|^2 |\theta_f|^2 C M_0^2 e^{-2|k|z_t} (1 - e^{-|k|(Z_b - Z_t)})^2 S^2(a, b) \quad (6)$$

where F is Fourier power spectrum, k is wavenumber (cycles km^{-1}), C_m is a constant, θ_m is a magnetization direction factor, θ_f is a magnetic field direction factor, M_0 is the magnetization, Z_b is depth to the bottom of the magnetic sources, Z_t is depth to the top of the ensemble of magnetic sources, and $S^2(a, b)$ is the horizontal dimensions of sources factor.

Figure 4 reveals that the azimuthally averaged log power spectrum, which was calculated using the Geosoft 8.4 software (Geosoft Oasis Montaj, 2015), has two segments reflecting the deeper and shallower sources of the magnetic field. The depth to

each zone was calculated from the slope of each segment of the spectrum using **Eq. 7**:

$$h = -\frac{s}{4\pi} \quad (7)$$

where h is the depth, and s is the slope of the log power (energy) spectrum.

The depth of the deeper sources with wavenumbers between 0.0 and 0.0001 cycle km^{-1} (**Figure 4**) was calculated at 7.8 km below the flight height, whereas the depth to the shallower sources with wavenumbers between 0.0001 and 0.0007 cycle km^{-1} was calculated at 1.3 km below the flight height (i.e., 305 m).

Centroid method

Bhattacharyya and Leu (1977) proposed a technique for calculating the centroid of rectangular parallelepiped sources. Tanaka et al. (1999) calculated the CPD map for East and Southeast Asia by dividing the region into subregional data over about 40,000 km^2 . Blakely (1988) divided the magnetic data map over the Nevada area into subregions with approximate areas of 14,400 km^2 and estimated the CPD of the state of Nevada. The CPD in Bulgaria was calculated by Trifonova et al. (2009) using six subregions with a 300-km edge. Hsieh et al. (2014) divided the integrated magnetic anomaly data of Taiwan into square subregions (250 × 250 km^2) and calculated

TABLE 1 | A sample of the calculated CPD with a 40 × 40 km² window, the estimated geothermal gradient, and heat flow. The data misfit using the RMSE is also shown.

Window No.	UTM		Root mean square error (RMSE)	Curie point depth (CPD) (km)	Geothermal gradient (°C/km)	Heat flow (mW/m ²)
	X	Y				
1	224,571.3	2,116,266	1.8	3.8	142.1	355.3
10	217,004.6	2,126,266	0.37	4	135	337.5
11	227,004.6	2,126,266	0.99	3.5	154.3	385.7
15	267,004.6	2,126,266	1.02	8.2	65.9	164.6
20	217,004.6	2,136,266	1.42	2.4	225	562.5
24	257,004.6	2,136,266	1.26	3.9	138.5	346.2
26	277,004.6	2,136,266	1.15	4.8	112.5	281.3
28	197,004.6	2,146,266	0.91	5	108	270
32	237,004.6	2,146,266	0.65	1.6	337.5	843.8
34	257,004.6	2,146,266	0.3	3.1	174.2	435.5
38	197,004.6	2,156,266	0.12	5.4	100	250
40	217,004.6	2,156,266	0.12	2.9	186.2	465.5
41	227,004.6	2,156,266	0.09	0.93	581.2	1453
43	247,004.6	2,156,266	0.08	2.6	207.7	519.2
47	287,004.6	2,156,266	0.25	4.1	131.7	329.3
49	207,004.6	2,166,266	0.34	3.6	150	375
53	247,004.6	2,166,266	0.43	2.5	216	540
57	287,004.6	2,166,266	0.67	5.8	93.1	232.8
59	207,004.6	2,176,266	0.88	2.9	186.2	465.5
63	247,004.6	2,176,266	0.59	3.7	145.9	364.9
67	287,004.6	2,176,266	0.58	7	77.1	192.9
71	227,004.6	2,186,266	0.75	3.7	145.9	364.9
75	267,004.6	2,186,266	0.69	6.5	83.1	207.7
81	227,004.6	2,196,266	0.73	5.1	105.9	264.7
83	247,004.6	2,196,266	0.54	4.9	110.2	275.5
86	277,004.6	2,196,266	0.35	8	67.5	168.8
88	197,004.6	2,206,266	0.27	7.8	69.2	173.1
91	227,004.6	2,206,266	0.31	5.7	94.7	236.8
92	237,004.6	2,206,266	0.43	4.5	120	300
94	257,004.6	2,206,266	0.44	5.4	100	250
99	207,004.6	2,216,266	0.5	5.1	105.9	264.7
101	227,004.6	2,216,266	0.86	4.9	110.2	275.5
103	247,004.6	2,216,266	0.8	7.4	73	182.4
105	267,004.6	2,216,266	0.79	7.7	70.1	175.3
108	197,004.6	2,226,266	0.34	6.3	85.7	214.3
111	227,004.6	2,226,266	0.02	5.1	105.9	264.7
113	247,004.6	2,226,266	0.25	12.4	43.5	108.8

the 2-D fast Fourier transform power spectrum for each region to estimate the CPD map. The Curie temperature depths in northern Italy (the Alps and the Po Plain) were calculated by Speranza et al. (2016) by creating windows with a 100-km edge and a 50% overlap between the aeromagnetic data and data from the Earth Magnetic Anomaly Grid two over the study area. Yang et al. (2017) calculated the CPD of Southeast Tibet via a spectral analysis of satellite data magnetic anomalies. Finally, Blakely (1995) presented the power spectral density of the total magnetic field [$\phi_{\Delta T}(\mathbf{k}_x, \mathbf{k}_y)$], as follows:

$$\phi_{\Delta T}(\mathbf{k}_x, \mathbf{k}_y) = \phi_M(\mathbf{k}_x, \mathbf{k}_y) * 4\pi^2 C_M^2 |\theta_f|^2 e^{-2|k|Z_t} (1 - e^{-2|k|(Z_b - Z_t)})^2 \quad (8)$$

Equation (8) was reduced to Eq. (9). Assuming that $M(\mathbf{x}; \mathbf{y})$ is the layer's magnetization and a random function of $\mathbf{x}, \mathbf{y}$, this indicates that the power-density spectra of the magnetization $\phi_M(\mathbf{k}_x, \mathbf{k}_y)$ is a constant. Then, the averaged power spectrum of $\phi(|\mathbf{k}|)$ can be written as:

$$\phi(|\mathbf{k}|) = A e^{-2|k|Z_t} (1 - e^{-2|k|(Z_b - Z_t)})^2 \quad (9)$$

where A is a constant. The depth to the top of the ensemble of magnetic sources (Z_t) was derived from the slope of the high wavenumber segment of a radially averaged power spectrum $\ln(P(k)^{1/2})$, as follows:

$$\ln(P(k)^{1/2}) = A - |k| Z_t \quad (10)$$

Additionally, the centroid depth (Z_0) could be estimated from the low wavenumber segment of the spectrum (Tanaka et al., 1999):

$$\ln(P(k)^{1/2}/k) = B - |k| Z_0 \quad (11)$$

where B is a constant. Finally, Z_b could be calculated from the following formula (Okubo et al., 1985):

$$Z_b = 2Z_0 - Z_t \quad (12)$$

According to Blakely (1996), the window size was suggested to be five times that of the CPD. Therefore, the current study applied a moving window to subset the data into

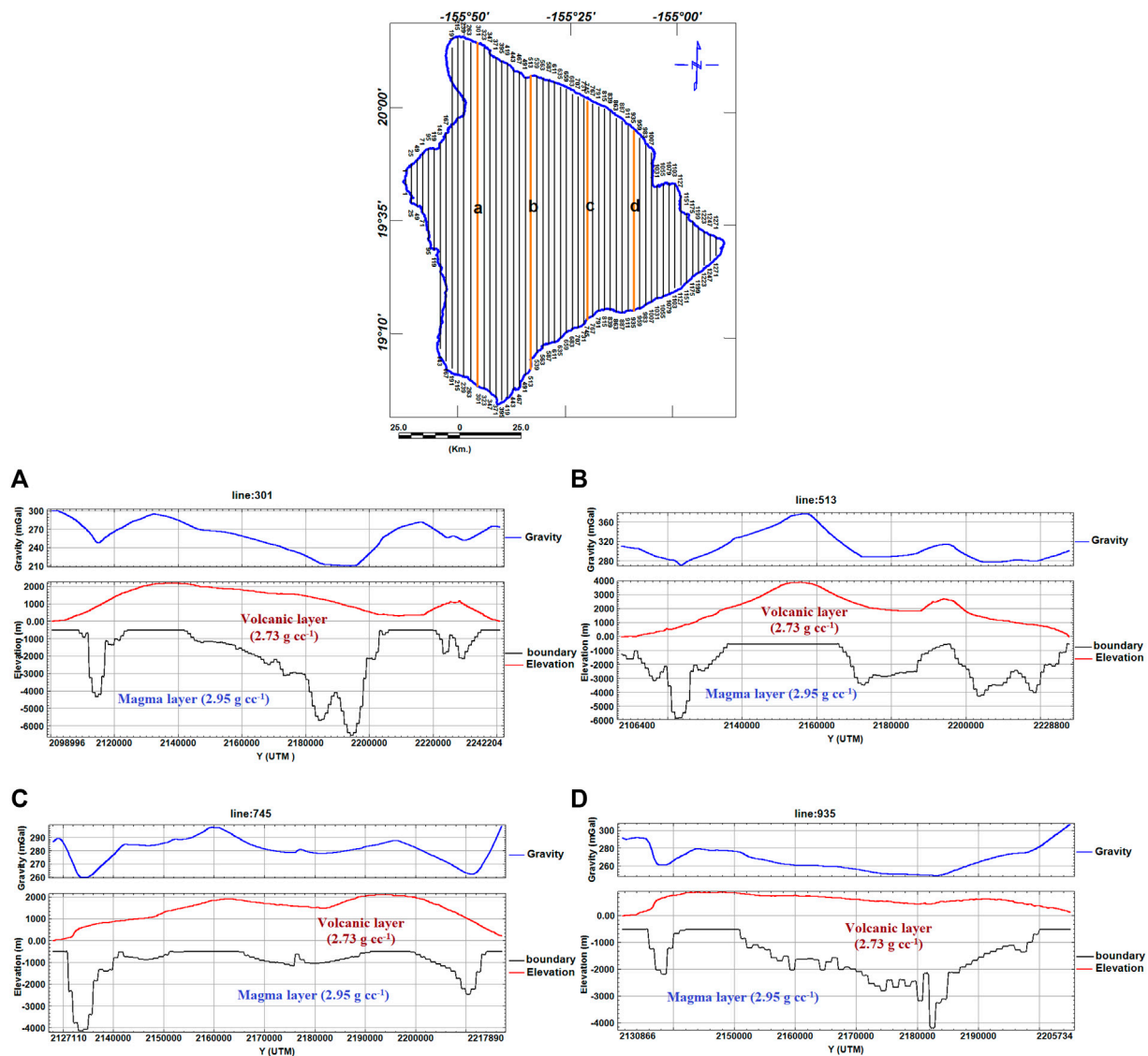


FIGURE 8 | Location map of the north–south parallel profiles used to calculate the density inversion model. The orange lines refer to selected 2-D modeled profiles (A–D).

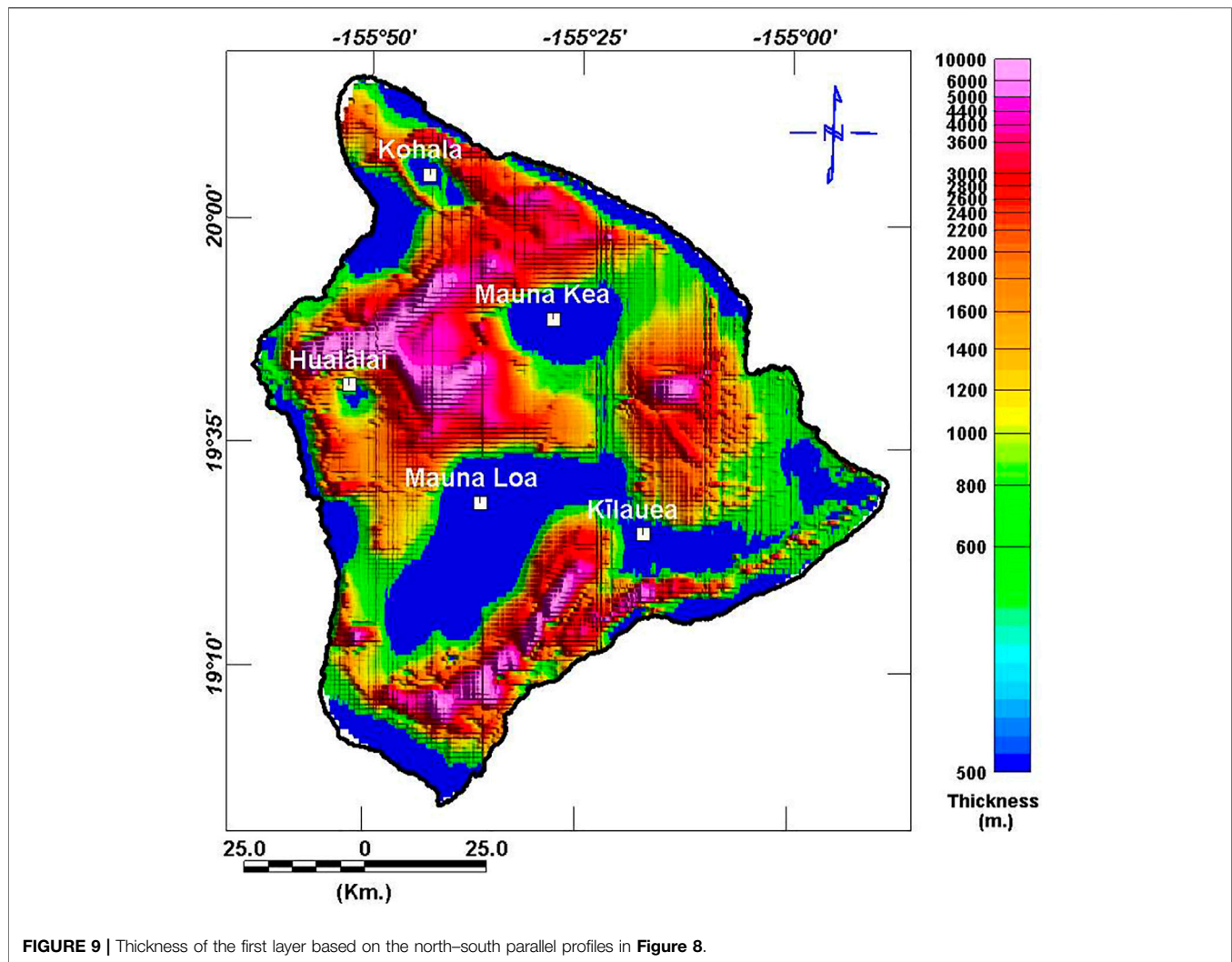
TABLE 2 | Pillow lava layers intercalated by hyaloclastite from the Hilo borehole and their densities [Source: (Moore, 2001)].

Lithologic zone	Depth range (m.)	Density (g/cc)
Hyaloclastite	Intercalated	2.3 to 2.7 Avg. 2.5
1 st pillow lava	1,983–2,136	3.01 ± 0.10
2 nd pillow lava	2,234–2,470	2.67 ± 0.13
3 rd pillow lava	2,640–2,790	2.89 ± 0.17
4 th pillow lava	2,918–3,097	2.97 ± 0.08

square regions with a 40-km edge and an overlap of 25% to generate adequate wavelengths for the power spectrum calculation.

An example of the azimuthally averaged power spectrum for some selected windows is given in **Figure 5** while

Figure 6 contains a location map of each spectral window. **Figure 7** shows the estimated depth to the bottom of the magnetized layer, which varies from 0.93 to 12.4 km below the ground surface when using a $40 \times 40 \text{ km}^2$ window. Based on the window of $40 \times 40 \text{ km}^2$, the surface of the magmatic plume was more correlated to the distribution of volcanoes on the surface of the island. Moreover, a large plume was observed beneath the ML, Ki, and MK volcanoes, underneath which uprising magma may be accumulating. **Table 1** presents the calculated CPD (using a $40 \times 40 \text{ km}^2$ window), the calculated geothermal gradient, and the heat flow. The misfit was also calculated (**Table 1**) using the root mean square errors (RMSEs) of the match between the observed and modeled spectra.



Morphology of the Magmatic Chamber from the Gravity Data

To construct a more realistic model for Hawaii Island, this section used the 2-D technique of gravity field modeling. The objective was to propose a comprehensive density model of the crust below the island.

Gravity inversion

The gravity data modeling was carried out along profiles (**Figure 8**) on the complete Bouguer anomaly map using the Intrepid 4.5 geophysics software (Intrepid Geophysics, 2013). This calculation method applies the algorithm presented by Murthy and Rao (1993). **Equation 13** provides the gravity anomaly $\Delta g(0)$ at any point $P(0)$ of a 2-D body of the polygonal cross section:

$$\Delta g(0) = 2Gd_c \sum_{k=1}^N p_k \left[\sin i_k \ln \left(r_k + \frac{1}{r_k} \right) - \cos i_k (\theta_{k+1} - \theta_k) \right] \quad (13)$$

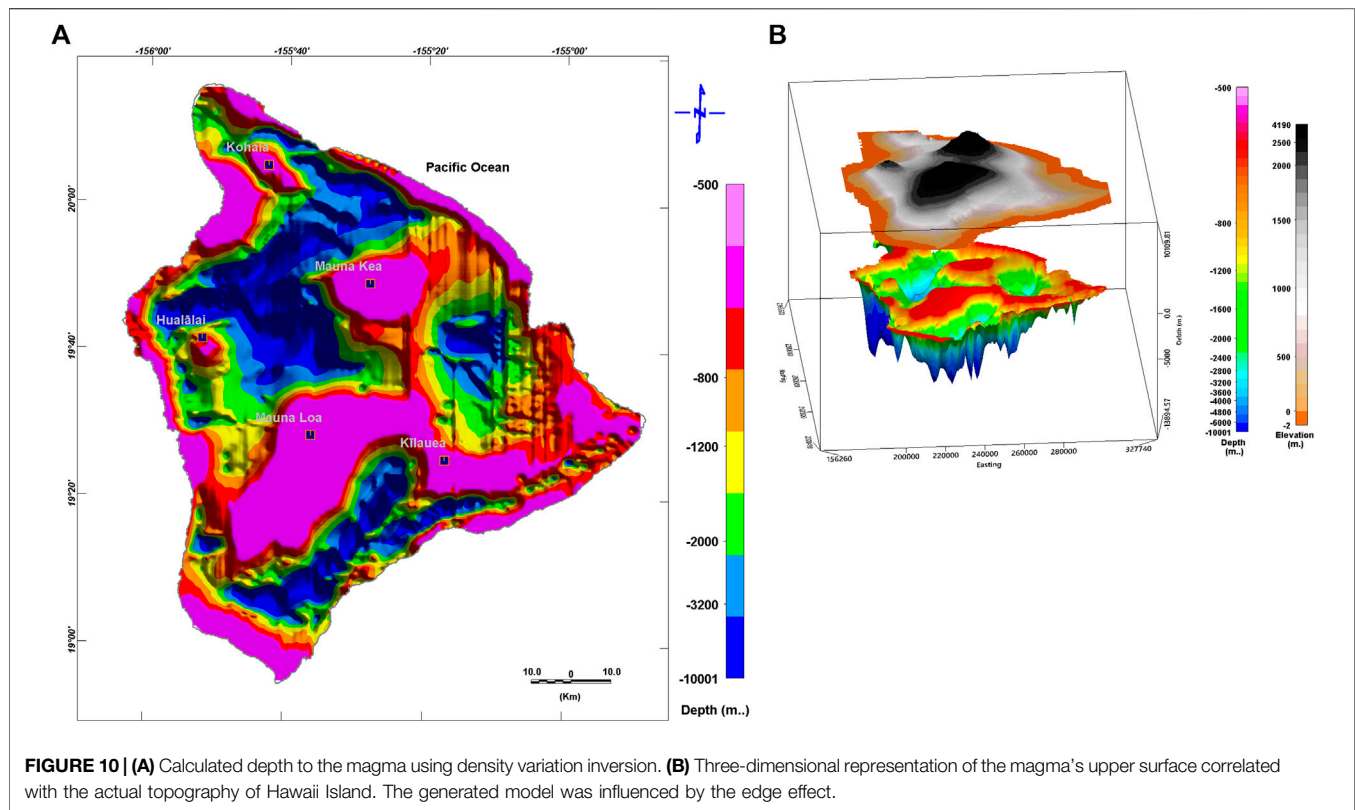
where G is the universal gravitational constant, r_k is the length of the line from the surface to the first point (vertex), r_{k+1} is the length to the next point, and θ_k and θ_{k+1} are the angles to the body vertices. The term r_k is defined by **Eq. 14**:

$$r_k = \sqrt{x_k^2 + z_k^2}, \quad r_{k+1} = \sqrt{x_{k+1}^2 + z_{k+1}^2} \quad (14)$$

where z_k is the depth to vertex k , and z_{k+1} is the depth to the next vertex of the assumed polygonal body.

This study constructed a 2-D density model for two Earth layers by forward modeling the gravity data along selected profiles (**Figure 8**). The two-layer model was constructed based on the density contrast properties of these layers and the assumption that the second layer (magma) was denser. The interface between the two layers was assumed to comprise a sequence of polygonal bodies, and a density contrast value of 0.22 g cc^{-1} between the two layers was also assumed.

Two-dimensional density models were operated along the profiles (with 1-km line spacing) using the GRAVIN module in the Intrepid 4.5 geophysics software (Intrepid Geophysics,



2013). Each model contained two layers. The first (top) volcanic layer had a density of 2.73 g cc^{-1} , which was averaged for 1,600 core samples obtained at depths between 889 and 3,097 m below sea level near the city of Hilo (Moore, 2001). The densities of the borehole samples are summarized in **Table 2**. The second magmatic layer was assumed to have a high density of 2.95 g cc^{-1} . As shown in **Figures 9, 10**, a complete two-layer solution for the gravity modeling was provided along the north–south profiles, which were interpolated across the island of Hawaii. **Figure 9** shows the thickness of the first layer based on the north–south parallel profiles in **Figure 8**, which equals the depth of the second (magma) layer in **Figure 10**. **Figure 10** depicts very steep vertical rising magmatic plumes, which resemble those calculated by the magnetic method. The Z_t value varied from 500 m to over 10 km. Compared with the actual volcanic activity conditions on Hawaii Island, these calculated depths were very reasonable. However, the edge effect meant that the outer rim of the model could not be considered.

GEOHERMAL GRADIENT AND HEAT FLOW

The geothermal gradient (dt/dz) from the Earth's surface to the CPD was estimated using **Eq. (15)** (Tanaka et al., 1999):

$$dt/dz = 540^\circ\text{C}/Z_b \quad (15)$$

The heat flow (q) was calculated using **Eq. 16** (Okubo et al., 1985):

$$q = h (dt/dz) = h(540^\circ\text{C}/Z_b) \quad (16)$$

This study used $h = 2.5 \text{ W/m}^\circ\text{C}$ as the value for igneous rocks (Springer, 1999). As can be seen in **Table 1**, the dt/dz varied from 43.5 to $581.1^\circ\text{C km}^{-1}$. **Figure 11** presents the estimated heat flow, which ranges between 108.8 and $1,453 \text{ mW/m}^2$ close to the volcanic eruption zone of ML. As geological constraints strongly influence CPDs (Ross et al., 2006), in the volcanic zones, the CPDs are shallower than 10 km (Obande et al., 2014). According to **Table 1** and **Figure 11**, Hawaii Island has high geothermal potential energy proximal to the volcanic eruption zones with shallower CPDs as well as high geothermal gradient and heat flows. The data in **Table 1** show window no. 213 to have the lowest heat flow, while the flow is highest in window no. 141.

DISCUSSION AND CONCLUSION

Investigating the morphology of the magmatic chamber beneath Hawaii Island as an example of an island arc system enhances the current understanding of plate tectonics theory. However, the growth and movement directions of these magmatic diapirs require a time-lapse investigation. Herein, we utilized the available potential data to better understand the morphology of the magmatic chamber, estimate the CPD, and calculate the heat flow.

This study successfully applied airborne geophysical data to characterize the morphology of the magmatic chamber beneath

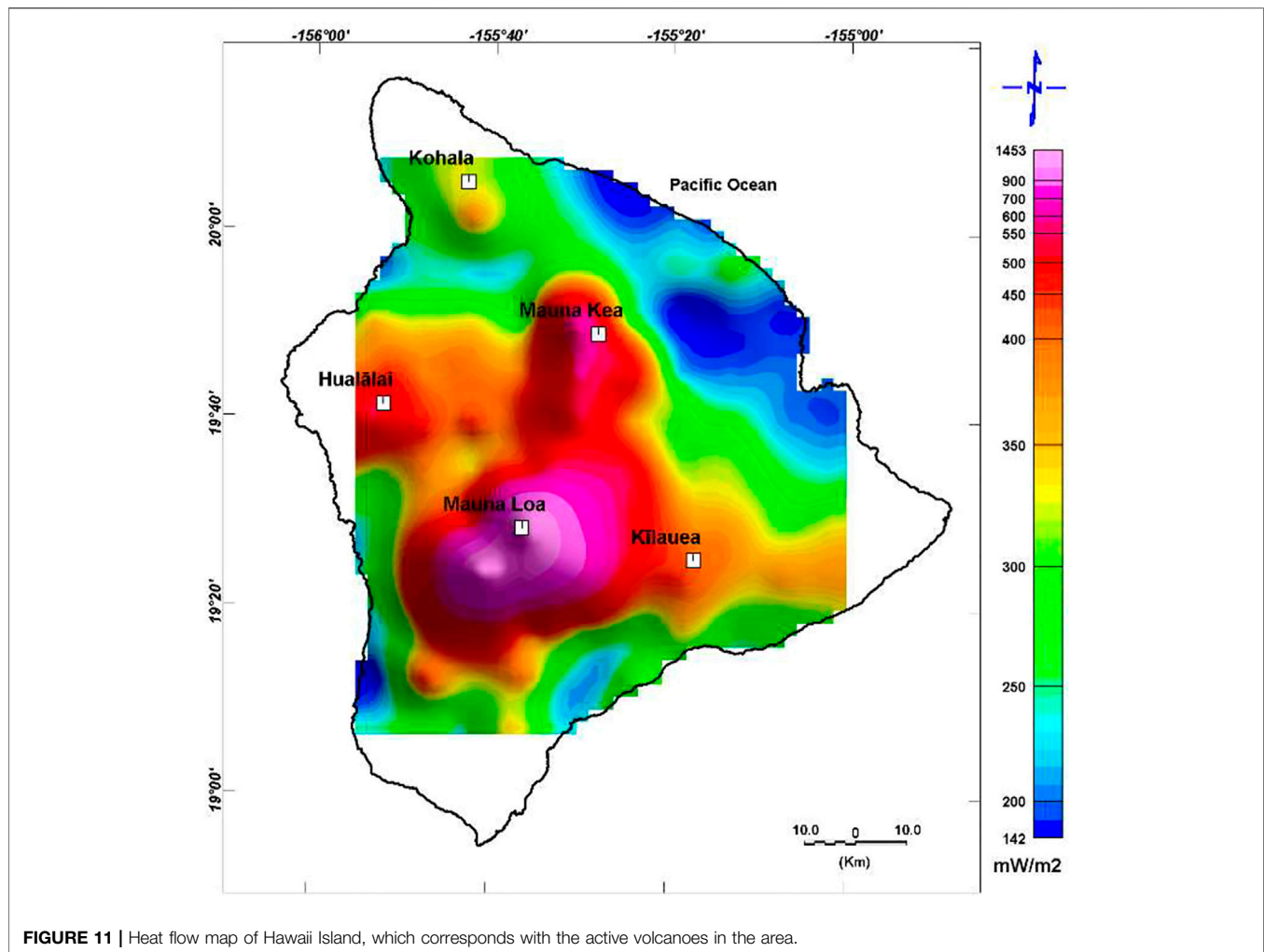


FIGURE 11 | Heat flow map of Hawaii Island, which corresponds with the active volcanoes in the area.

Hawaii Island and estimate the depths to the deeper sources; these varied from 7.8 km below the flight height to shallower sources estimated to be 1.3 km deep. The Curie point is located from shallow depths (<1 km) at and proximal to the ML volcano to 12 km close to the northeastern part of the MK volcano region below the ground surface.

The results of the gravity data inversion revealed the existence of very steep vertical rising magmatic plumes, reflecting the dense intrusive material of basaltic and ultramafic rocks underlying the volcanoes. The results corresponded well with those calculated using the magnetic method. The depth to the top of the second (magma) layer varied from 500 m to about 10 km. An inverse relation was identified between the heat flow and the CPD, whereby a decrease in heat flow corresponded with an increase in the CPD. The highest value of 1,453 mW/m² occurred at the shallowest CPD (0.93 km) in window no. 141, while the lowest value of 108.8 mW/m² was correlated with the highest CPD (12.4 km) in window no. 213. According to Jessop et al. (1976), these anomalous geothermal conditions are related to heat flows of between 80 and 100 mW/m².

The results of this study are in good agreement with the previous 3-D gravity models used by Kauahikaua et al. (2000),

who studied the magmatic structures within basaltic volcanoes and defined the structures related to seismic hazards and landslides on the island of Hawaii. Their results revealed the existence of dense cumulates and intrusions beneath the summits of every volcano.

Information on the magmatic properties of subsurface rocks was provided by Hildenbrand et al. (1993), who applied the spectral depth technique to aeromagnetic data. Two shallow magnetic zones were identified at a depth of 1 km. Additionally, using the spectral inversion of the magnetic data, they defined a deep magnetic horizon at 10.5 km below the flight height. The results of the current study correspond well with those of Hildenbrand et al. (1993), as our interpretation revealed shallow and deep magnetic sources to be located at 1.3 and 8 km below the flight height, respectively. Based on the CPD estimations, the bottom of the magnetized layer varied between 0.93 and 12.4 km in depth, while the gravity inversion technique indicated a depth of 0.5–10 km below the ground surface.

The existence of large magmatic plumes of dense cumulates and intrusions beneath the summits of the MK, ML, and Ki active volcanoes in this study was confirmed by the presence of

higher-velocity regions near the active volcanoes containing dense intrusive materials, such as olivine cumulates (Thurber, 1984; Hill and Zucca, 1987; Okubo et al., 1997). However, a 3-D P-wave velocity model run by Park et al. (2007) for the southeastern part of Hawaii Island produced higher crustal seismic velocity values of 7.0–7.4 km/s. These higher values indicate the dominance of dense olivine cumulates mixed with extrusive and intrusive basaltic rocks in the high-velocity regions underlying the summits and rift regions of the active areas of Ki and ML. In the study of Park et al. (2007), the lower-velocity regions south of the Hilina and along the Kao'iki fault zones were attributed to thick accumulations of volcanoclastic sediments.

Using more than one geophysical dataset provides adequate confidence in the calculated results to overcome the problem of nonuniqueness. The calculated upper surface of the magma chamber delineated about five diapirs, three of which (ML, MK, and Ki) were grouped in one large plume of dense basic and ultrabasic materials. Compared with other studies, the integrated approach presented herein provides an in-depth understanding of the geometry of the magmatic chamber using an indirect CPD estimation to calculate the depth to the chamber's upper surface. Finally, the new gravity model for Hawaii Island, which incorporates advanced software, corresponds well with the results of the magnetic data.

KEY POINTS

Estimates the depth to the bottom of the magnetized rocks based on calculating the Curie point depth beneath Hawaii Island.

Characterizes the morphology of the magmatic chamber beneath Hawaii Island using the gravity inversion technique.

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- Estimates the geothermal gradient and heat flow associated with the active volcanoes of Hawaii Island.
- DATA AVAILABILITY STATEMENT**
- The aeromagnetic field data (Hawaii-78-Hawaii, Godson et al., 1981) can be found at http://mrdata.usgs.gov/geophysics/surveys/geophysics2/HI/HI_1071.jpg and http://mrdata.usgs.gov/geophysics/surveys/geophysics2/HI/HI_1071.zip. The gravity measurements for Hawaii Island (Kauahikaua, 2017) can be found at <https://doi.org/10.5066/F7V1230Q>.
- AUTHOR CONTRIBUTIONS**
- All authors listed have made a substantial, direct, and intellectual contribution to the work and approved it for publication.
- ACKNOWLEDGMENTS**
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